Jet discrimination with a quantum complete graph neural network Yi-An Chen, Kai-Feng Chen

Department of Physics, National Taiwan University, Taipei, Taiwan

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Outline

- **Classical Machine Learning**
 - Message-Passing Graph Neural Network
- **Quantum Machine Learning** •
 - Variational Quantum Circuit
 - Quantum Complete Graph Neural Network •
- **Training X IBMQ X Summary**
 - Public Monte Carlo Simulated Jet Data lacksquare
 - Running on IBMQ



New Quantum Algorithm

Variational Quantum Algorithm



Classical Neural Networks **Models for Jet Discrimination**

Machine Learning

Analysis Done

Data





Classical ML for Jets Overview



DNN (arXiv 1704.02124)



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LSTM (arXiv 1711.09059)



Classical Machine Learning Message-Passing Graph Neural Network (MPGNN)

- Message Passing: Compute the information for each particle pair through some parametrized transformation.
- 2. Node Aggregation: Aggregate the transformed information for each particle. Typically elementwise summation \iff permutation-invariant
- 3. Repeat step 1 & 2 for several times (optional).
- Graph Aggregation: Aggregate the information of 4. all particles.

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3. Repeat steps 1 & 2

B

4. Graph Aggregation

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Classical Machine Learning Deep Sets Theorem and MPGNN

usually written as:

$$\mathbf{x}_{i}^{(k)} = \gamma^{(k)} \left(\mathbf{x}_{i}^{(k-1)}, \bigoplus_{j \in \mathcal{N}(i)}^{\mathbf{k}} \right)$$

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Deep Sets Theorem (arXiv 1703.06114) : A function (model) f is permutation-invariant over a

set *X* (particles) if and only if $f(X) = g\left(\sum_{\mathbf{x}_i \in X} h(\mathbf{x}_i)\right)$ for some suitable transformations *g* and *h*.

The Message-Passing Graph Neural Network (MPGNN) obeys the Deep Sets Theorem, and is

Aggregation function (MEAN, SUM, MAX, etc.)

QCGNN Quantum Complete Graph Neural Network

QCGNN **Quantum Complete Graph Neural Network**

Suppose we have N particles with features $\{\mathbf{x}_i \mid 0 \le i \le N-1\}$. We prepare a quantum circuit with $n_I + n_O$ qubits where

- $n_I = \lceil \log_2 N \rceil$ is the number of qubits in the index register (*IR*)
- n_O is the number of qubits in the network register (NR)

The initial quantum state is initialized as

$$|\psi_0\rangle = \frac{1}{\sqrt{N}} \sum_{i=0}^{N-1} |i\rangle|0\rangle$$

If $N = 2^{n_I}$, then we can simply use Hadamard gates. Otherwise, one should use some Uniform State Oracle (USO) to prepare the state.

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 $\otimes n_O$

Uniform State Oracle

QCGNN **Quantum Circuit**

QCGNN Measurement

The quantum state before measurement is $|\psi\rangle = \frac{1}{\sqrt{N}} \sum_{i=0}^{N-1} |\mathbf{x}_i, \theta\rangle$

Consider a Hermitian matrix J with dimension $2^{n_I} \times 2^{n_I}$ full of ones, i.e., ~

X

 $J = \begin{bmatrix} 1 & \dots & 1 \\ \vdots & & \vdots \\ 1 & & 1 \end{bmatrix} \stackrel{\checkmark}{=} (I + X)^{\bigotimes n_I}.$

Denote the observables on NR as P, then, $\langle \psi | J \otimes P | \psi \rangle = \frac{1}{N} \sum_{i < N} \sum_{j < N} \langle \mathbf{x}_i; \theta | P | \mathbf{x}_j; \theta \rangle$

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Similar to MPGNN, with automatic aggregation.

Experiments Results Quantum Simulators & IBMQ

Jet Dataset

- Top-Taggers (arXiv 1902.09914)
 - Top v.s. Gluon / Light Quarks
 - Jet $p_T \in [550, 650] GeV$
 - Pythia / Delphes (ATLAS)
 - FastJet with R = 0.8
- JetNet (arXiv 2106.11535)
 - Multi-class $\{g, q, t, W, Z\}$
 - Jet $p_T \sim 1 TeV$
 - MadGraph / Pythia
 - FastJet with R = 0.8

events of Number

ents 150K e < 100K of Number

Only the 30 highest p_T particles are provided March 18, 2025 ISGC @ Academia Sinica

Training Results AUC and Accuracy

- Each training process was conducted with 5 different random seeds and 30 epochs.
- Each class has 25K training samples, 2.5K validation samples, and 2.5K testing samples.
- The number of particles of jets lies between $4 \sim 16 \implies$ At most $n_I = 4$ qubits for IR is needed.
- The performance of the state-of-the-art classical models is also presented.

State-of-the-art classical models		Madal	TOP Dataset (2 classes)			JETNET Dataset (5 classes)		
		Wiodei	# params	• AUC	Accuracy	# params	AUC	Accura
	$\langle ($	Particle Transformer	$2.2\mathrm{M}$	$0.946 {\pm} 0.005$	$0.868 {\pm} 0.009$	$2.2\mathrm{M}$	$0.889 {\pm} 0.002$	0.656 ± 0.656
		Particle Net	$177 \mathrm{K}$	$0.953{\pm}0.003$	$0.885{\pm}0.006$	$178\mathrm{K}$	$0.896{\pm}0.003$	0.669 ± 0.669
Classical models for		Particle Flow Network	$72.3\mathrm{K}$	$0.954{\pm}0.004$	$0.885{\pm}0.005$	$72.7\mathrm{K}$	$0.900 {\pm} 0.003$	0.675 ± 0.675
		$MPGNN - n_M = 64$	$13\mathrm{K}$	$0.961{\pm}0.003$	$0.896{\pm}0.003$	$13.3\mathrm{K}$	$0.903 {\pm} 0.002$	0.683 ± 0.683
Denchmarking	\geq	$MPGNN - n_M = 6$	255	$0.924{\pm}0.006$	$0.866 {\pm} 0.006$	323	0.865 ± 0.004	$0.615 \pm 0.$
		MPGNN - $n_M = 3$	126	$0.922{\pm}0.005$	$0.864{\pm}0.006$	194	$0.757 {\pm} 0.110$	0.475 ± 0.00
OCGNN with		QCGNN - $n_Q = 6$	201	$0.932{\pm}0.004$	$0.868 {\pm} 0.005$	269	$0.822{\pm}0.003$	0.543 ± 0.543
v = 3 and $v = 6$		QCGNN - $n_Q = 3$	99	$0.919 {\pm} 0.006$	$0.864{\pm}0.005$	167	$0.796{\pm}0.009$	0.505 ± 0.505
$n_Q = 3$ and $n_Q = 6$								
(On simulators)								
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IBMQ Results Runtime of Quantum Gates

- T_{ENC} and T_{PARAM} are the time for encoding and parametrized gates respectively.

IBMQ Backend	Ν	$T_{\rm ENG}$
	2	2.56'
ibm_nazca	4	5.352
	8	10.55
	• 2 •	2.59!
2 2.598 bm_strasbourg 4 5.416		
	8	11.08

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The gate runtime experiment is conducted with two different IBMQ backends for 10 times.

	a de la companya de la		
C		$T_{\rm PARAM}$	$\Box = Scales as O(N)$
7		0.209	
2		0.197	
51		0.219	
č		0.217	
6	2	0.197	
5		0.211	Constant time $O(1)$

Summary

- In the task of jet discrimination, graph is one of the natural representation. To design
- classical MPGNN requires $O(N^2)$.
- 0 quantum circuits, information transmission was unsuccessful.
- advantage of the QCGNN can be studied.

permutation-invariant models, graph neural networks have become a popular architecture.

• We introduce a new quantum model, the Quantum Complete Graph Neural Network. If the parametrized gates are deep enough, the cost of QCGNN only scales as O(N), while

QCGNN has also been tested on IBMQ real quantum devices. However, due to noise in the

• As the quantum computers becoming more robust in the future, the potential for quantum

Backup Slides

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CMS Coordinate System

The particle flow is defined as

• The transverse momentum $p_T = \sqrt{p_x^2 + p_y^2}$

• The azimuthal angle $\phi = \tan^{-1}(\frac{p_y}{n})$

• The pseudo-rapidity $\eta = -\ln\left[\tan(\frac{\theta}{2})\right]$

In jet analysis, the differences $\Delta \phi$ and $\Delta \eta$ relative under boosts in *z*-direction.

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CMS Coordinate System by Izaak Neutelings

In jet analysis, the differences $\Delta \phi$ and $\Delta \eta$ relative to the jets is adopted, since they are Lorentz invariant

Dataset **Top Tagging (arXiv 1902.09914)**

Data set 2

For the signal only, we further require a matched parton-level top to be within $\Delta R = 0.8$, and all top decay partons to be within $\Delta R = 0.8$ of the jet axis as well. No matching is performed for the QCD jets. We also require the jet to have $|\eta_j| < 2$. The constituents are extracted through the Delphes energy-flow algorithm, and the 4-momenta of the leading 200 constituents are stored. For jets with less than 200 constituents we simply add zero-vectors.

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The top signal and mixed quark-gluon background jets are produced with using Pythia8 [25] with its default tune for a center-of-mass energy of 14 TeV and ignoring multiple interactions and pile-up. For a simplified detector simulation we use Delphes [26] with the default ATLAS detector card. This accounts for the curved trajectory of the charged particles, assuming a magnetic field of 2 T and a radius of 1.15 m as well as how the tracking efficiency and momentum smearing changes with η . The fat jet is then defined through the anti- k_T algorithm [27] in FastJet [28] with R = 0.8. We only consider the leading jet in each event and require

$p_{T,j} = 550 \dots 650 \text{ GeV}$.

(1)

Dataset JetNet (arXiv 2106.11535)

JetNet Generation B

LHC, so as to remain experiment-independent and allow public access to the dataset.

The so-called parton-level events are first produced at leading-order using MAD-GRAPH5_aMCATNLO 2.3.1 [51] with the NNPDF 2.3LO1 parton distribution functions [52]. To focus on a relatively narrow kinematic range, the transverse momenta of the partons and undecayed gauge bosons are generated in a window with energy spread given by $\Delta p_{\rm T}/p_{\rm T} = 0.01$, centered at 1 TeV. These parton-level events are then decayed and showered in PYTHIA 8.212 [5] with the Monash 2013 tune [53], including the contribution from the underlying event. For each original particle type, 200,000 events are generated. Jets are clustered using the anti- $k_{\rm T}$ algorithm [54], with a distance parameter of R = 0.8 using the FASTJET 3.1.3 and FASTJET CONTRIB 1.027 packages [55, 56]. Even though the parton-level $p_{\rm T}$ distribution is narrow, the jet $p_{\rm T}$ spectrum is significantly broadened by kinematic recoil from the parton shower and energy migration in and out of the jet cone. We apply a restriction on the measured jet $p_{\rm T}$ to remove extreme events outside of a window of $0.8 \text{ TeV} < p_T < 1.6 \text{ TeV}$ for the $p_T = 1 \text{ TeV}$ bin. This generation is a significantly simplified version of the official simulation and reconstruction steps used for real detectors at the

Particle Flow Network arXiv 1810.05165

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Motivated by the Deep Set Theorem

Μ $\mathcal{O}(\{p_1,\ldots,p_M\})=F$ $\Phi(p_i)$

Single Qubit Gates **Rotation Gates**

$$R_{x}(\theta) \equiv e^{-i\theta X/2} = \cos\frac{\theta}{2}I - i\sin\frac{\theta}{2}X = \begin{bmatrix} \cos\frac{\theta}{2} & -i\sin\frac{\theta}{2} \\ -i\sin\frac{\theta}{2} & \cos\frac{\theta}{2} \end{bmatrix}$$
(4.4)

$$R_{y}(\theta) \equiv e^{-i\theta Y/2} = \cos\frac{\theta}{2}I - i\sin\frac{\theta}{2}Y = \begin{bmatrix} \cos\frac{\theta}{2} & -\sin\frac{\theta}{2} \\ \sin\frac{\theta}{2} & \cos\frac{\theta}{2} \end{bmatrix}$$
(4.5)

$$R_{z}(\theta) \equiv e^{-i\theta Z/2} = \cos\frac{\theta}{2}I - i\sin\frac{\theta}{2}Z = \begin{bmatrix} e^{-i\theta/2} & 0 \\ 0 & e^{i\theta/2} \end{bmatrix} .$$
(4.6)

Theorem 4.1: (Z-Y decomposition for a single qubit) Suppose U is a unitary operation on a single qubit. Then there exist real numbers α, β, γ and δ such that

$U = e^{i\alpha} R_z(\beta) R_y(\gamma) R_z(\delta).$ (4.11)

Multi-Qubit Gates Decomposition of Multi-Controlled Gates

Figure 4.10. Network implementing the $C^n(U)$ operation, for the case n = 5.

Adopted from "Quantum Computation and Quantum Information" by N & C. March 18, 2025 ISGC @ Academia Sinica

Uniform State Oracle arXiv 2306.11747

In this paper, we propose an efficient approach for quantum state preparation of uniform superposition state $|\Psi\rangle = \frac{1}{\sqrt{M}} \sum_{j=0}^{M-1} |j\rangle$ that offers a significant (exponential) reduction in gate complexity and circuit depth without the use of ancillary qubits. We show that using only $n = \lceil \log_2 M \rceil$ qubits, the uniform superposition state $|\Psi\rangle$ can be prepared for arbitrary M with a gate complexity and circuit depth of $O(\log_2 M)$.

VQC Ansatz PennyLane

qml.BasicEntanglerLayers

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 $R(lpha_1^1,eta_1^1,\gamma_1^1)$ $R(lpha_2^1,eta_2^1,\gamma_2^1)$ $R(\alpha_3^1,\beta_3^1,\gamma_3^1)$ $R(lpha_4^1,eta_4^1,\gamma_4^1)$

qml.StronglyEntanglingLayers

Model Setup

Noise Simulated with PennyLane

Data Encoding	
Noise	
VQC	
Noise	
Data Encoding	
Noise	
VQC	
Noise	

Parameter Shift Rule arXiv 1905.13311

Gavin E. $Crooks^*$ California Institute of Technology, Pasadena, CA 91125, USA and Berkeley Institute for Theoretical Sciences, Berkeley, CA 94706, USA

- observable and $U_G(\theta) = e^{-ia\theta G}$ with some Hermitian operator G.
- If G has two unique eigenvalues e_0 and e_1 , the gradient can be calculated by

$$\frac{d}{dx}f(x) = r\left[f(\theta + \frac{\pi}{4r}) - f(\theta - \frac{\pi}{4r})\right] \quad \text{with} \quad r = \frac{a}{2}(e_1 - e_0)$$

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Gradients of parameterized quantum gates using the parameter-shift rule and gate decomposition

• Consider a VQC output $f(\theta) = \langle \psi | U_G^{\dagger}(\theta) A U_G(\theta) | \psi \rangle$, where A is some Hermitian operator of

Barren Plateau arXiv 2309.09342

